

LA-UR-14-27226

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Investigation of light-matter interactions: Photoluminescence properties of individual quantum emitters Title:

Author(s): Ma, Xuedan

Intended for: Progress meeting for a possible collaboration

2014-10-22 (Rev.1) (Draft) Issued:





Investigation of Light-Matter Interactions: Photoluminescence Properties of Individual Quantum Emitters

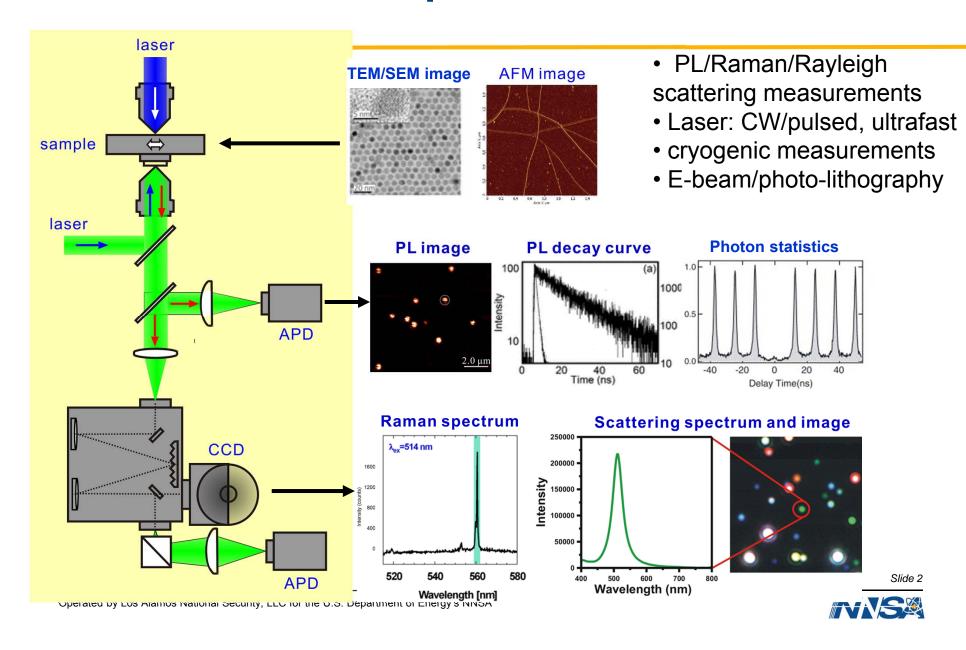
Xuedan Ma

Center for Integrated Nanotechnologies, Materials Physics & Applications Division
Los Alamos National Laboratory



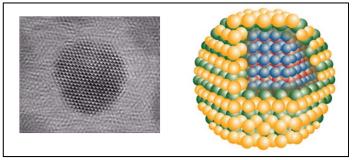


Research Techniques

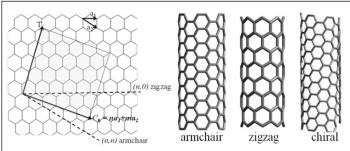


Mainly Studied Topics

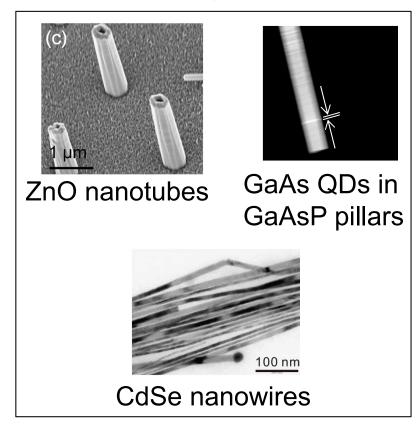
II-VI semicondutor quantum dots (QDs)



Single-walled carbon nanotubes (SWCNTs)



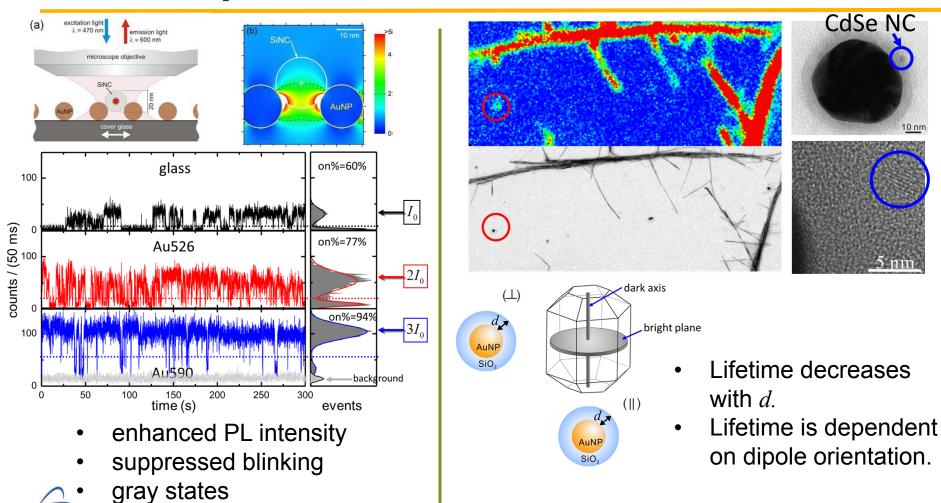
Semiconductor nanotubes/nanowires





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QDs Coupled to Surface Plasmons



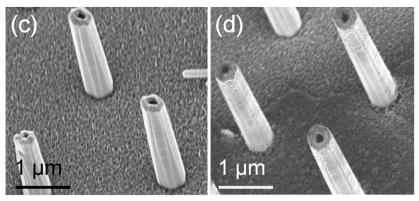
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Ma et al. Nano Lett. 2010, 10, 4166 Ma et al. J. Phys. Chem. Lett. 2011, 2, 2466 Ma et al. J. Phys. Chem. C 2013, 117, 16698

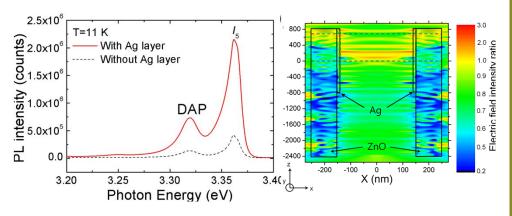


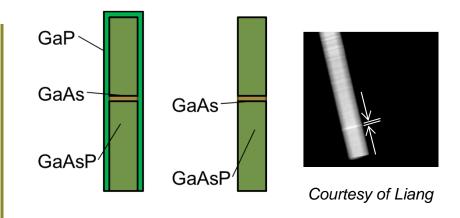


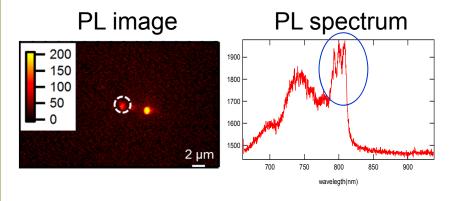
Semiconductor Nanotubes/Nanowires



ZnO NT ZnO NT + Pt (inner)









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Z. Li, X. Ma et al. Eur. J. Inorg. Chem. 2010, 27, 4325 Z. Wang, X. Ma et al. Small 2011, 7, 2464

J. Yoo, X. Ma et al. Nano Lett. 2013, 13, 2134





- Photon Statistics of Individual Carbon Nanotubes at Room Temperature
- Oxygen Doping Modifies Excitonic Fine Structures of Carbon Nanotubes
- Influence of Exciton Dimensionality on Spectral Diffusion of Single-Walled Carbon Nanotubes





Photon Statistics of Individual Carbon Nanotubes at Room Temperature

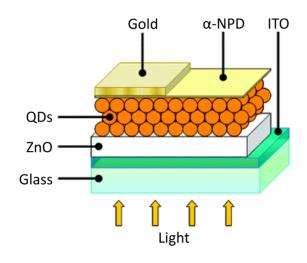


Ma et al. Phys. Rev. Lett. in submission



Why Photon Statistics?

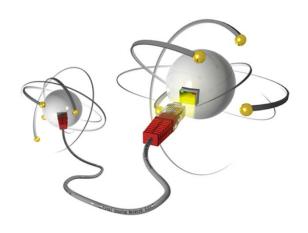
photovoltaics and solar cells



ACS Nano, 3, 3638-3648 (2009), Aydil et al.

multiexciton generation rate

quantum information



A. Neuzner, MPQ Nature, 484, 195 (2012), Ritter et al.

single photon source

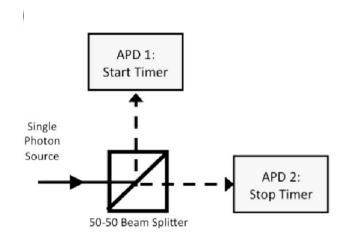


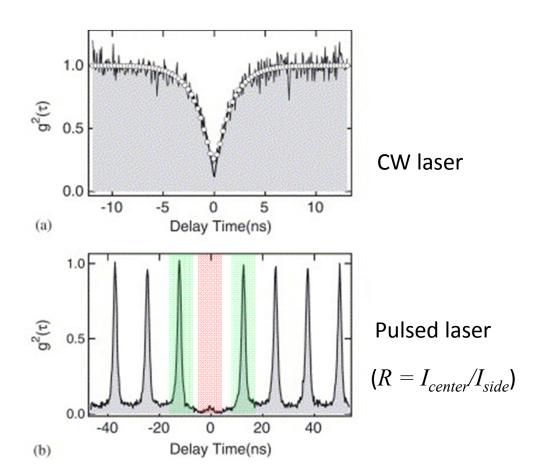
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Second-order Correlation Function g⁽²⁾:

Hanbury-Brown-Twiss Setup



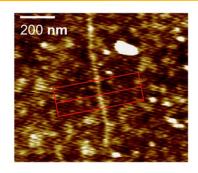


• g⁽²⁾: a measure of photon statistics.



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Photon Statistics of CNTs: State of the Art



4.2 K

Polymer 160 nm

Polymer 160 nm

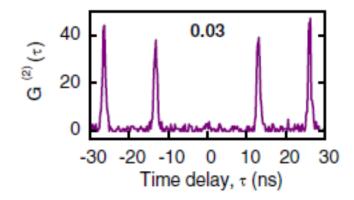
Au 100 nm

SiO2 100 nm

9 K

PRL 100, 217401 (2008), Hoegele et al.

Nano Lett. 12, 1934 (2012), Stefan et al.



R < 0.5photon antibunching

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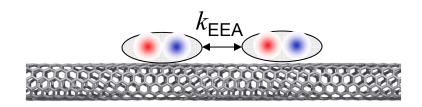
• Photon antibunching: photons are separated from each other.

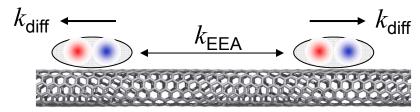
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Photon Statistics of CNTs: Mechanism

Low T

Room T





localized exciton:

exciton – exciton annihilation (EEA) diffusing exciton:

- exciton exciton annihilation (EEA)
- exciton diffusion



How exciton diffusion and EEA affect photon statistics in 1D system?

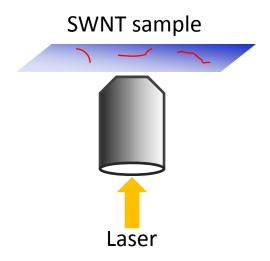


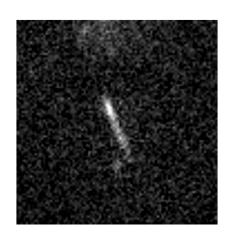
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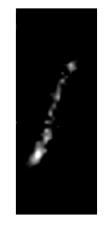


Individual SWNT Imaging

Wide field images







Stable PL

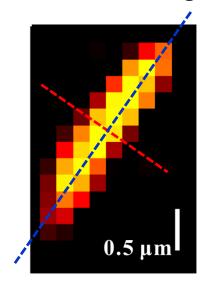
PL blinking/fluctuation

- Homogeneous illumination of the whole SWNTs with expanded laser beam
- Beam size: ~60 x 60 um
- SWNT length: <15 um

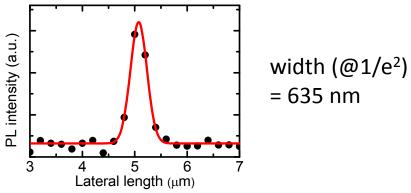


Individual SWNT Imaging

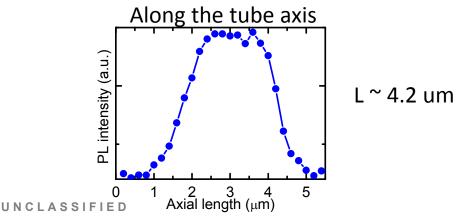
Wide field image



Perpendicular to the tube axis

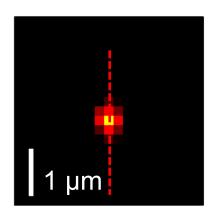


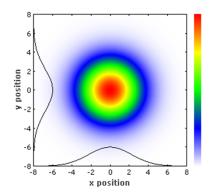
- Fitted with a Gaussian function
- Diffraction limited

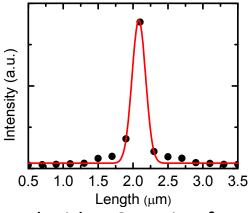




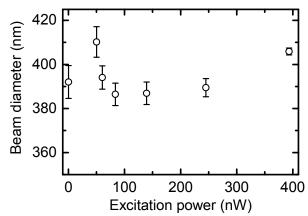
Gaussian Laser Spot







- Fitted with a Gaussian function
- Diffraction limited

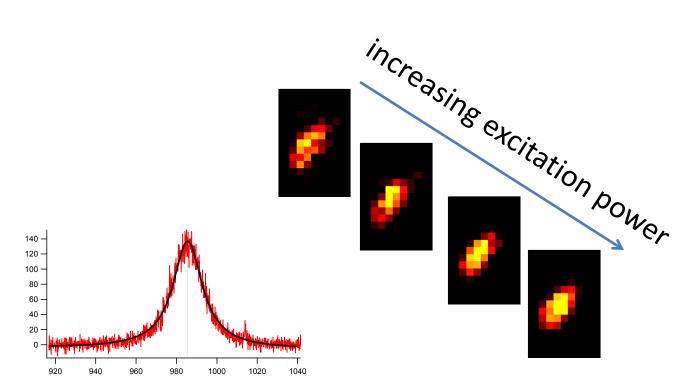


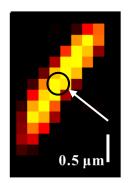
• Beam size independent of power.



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Excitation Power Dependent PL images





- 985.51 nm, 17.7 nm
- Fitted with a Lorentz function

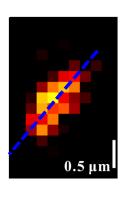
elongated PL image

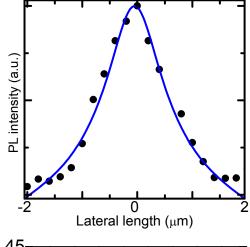


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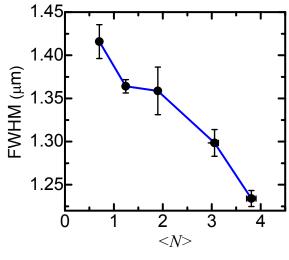


PL Intensity Profile





- width > point spread function
- due to exciton diffusion



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Profile width changes with excitation power.

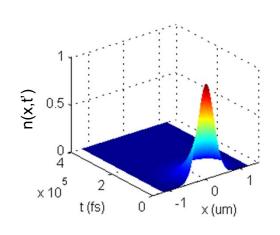


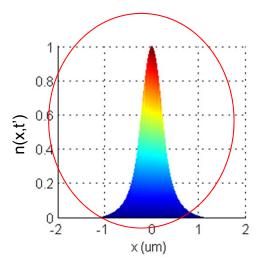
1D Diffusion Equation

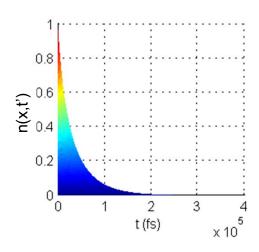
$$\frac{\partial n(x,t')}{\partial t'} = -n(x,t') + \underbrace{L_D^2}^{2} \underbrace{\frac{\partial^2 n(x,t')}{\partial x^2}}_{D} - \underbrace{C} \cdot n(x,t')^2$$

diffusion length

exciton-exciton annihilation coefficient







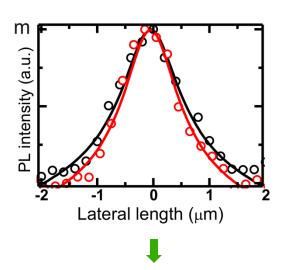
Intensity profile

decay curve

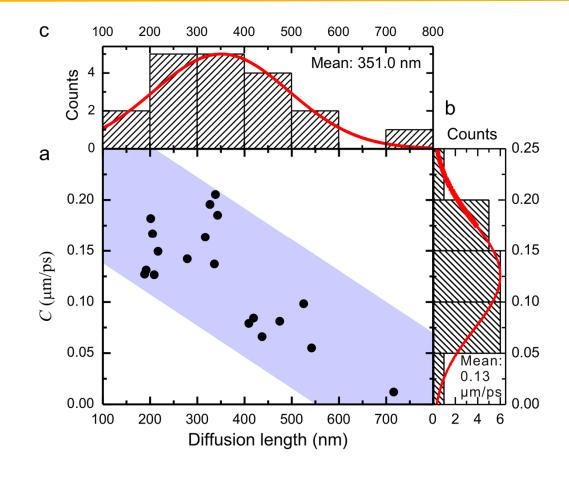


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Fitting Intensity Profiles



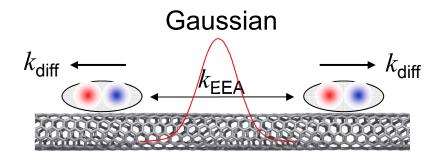
- diffusion length L_D
- EEA coefficient C





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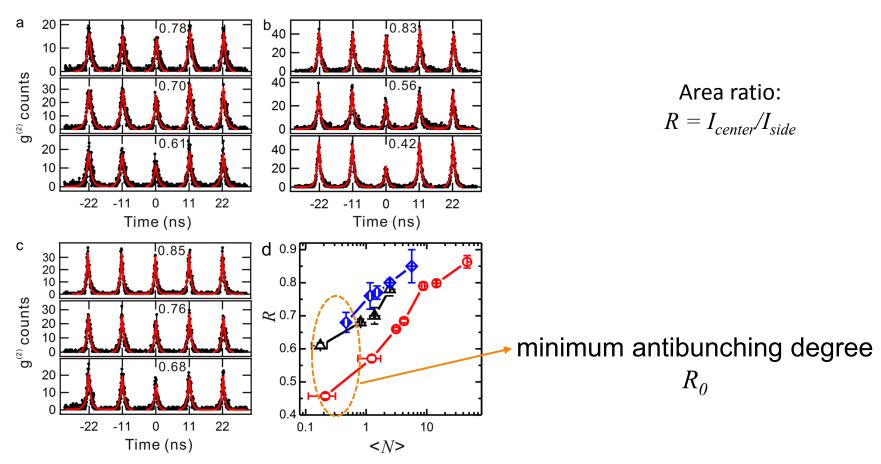
Exciton Diffusion Exciton-Exciton Annihilation



- Larger diffusion length/rate
 - → less efficient exciton-exciton annihilation
 - \rightarrow smaller C



Excitation Power Dependent g⁽²⁾ Measurements



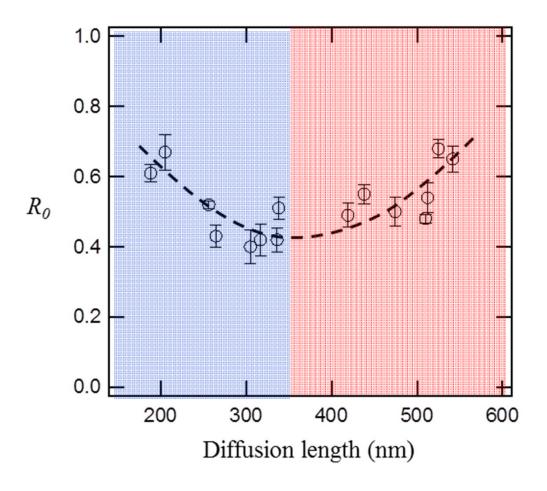
R increases with increasing excitation power.

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$R_{\theta} \Leftrightarrow$ diffusion length L_{D}





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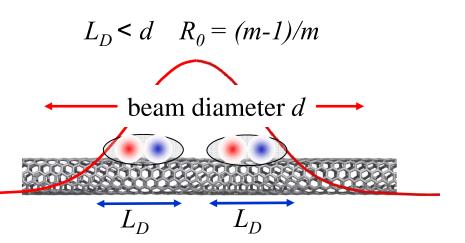


Minimum Antibunching Degree R_{θ}

$$R_0 = Q_{2X} / (mQ_{1X}) + (m-1) / m$$

 Q_{2X} : emission efficiency of two excitons Q_{1X} : emission efficiency of one exciton m: d/L_D

$$L_D > d \quad R_0 = Q_{2X}/(mQ_{1X})$$
+ beam diameter d +
$$-$$
diffusion length L_D



- Section of SWCNT behaviors like
 a QD.
- Excitons behave like independent emitters.

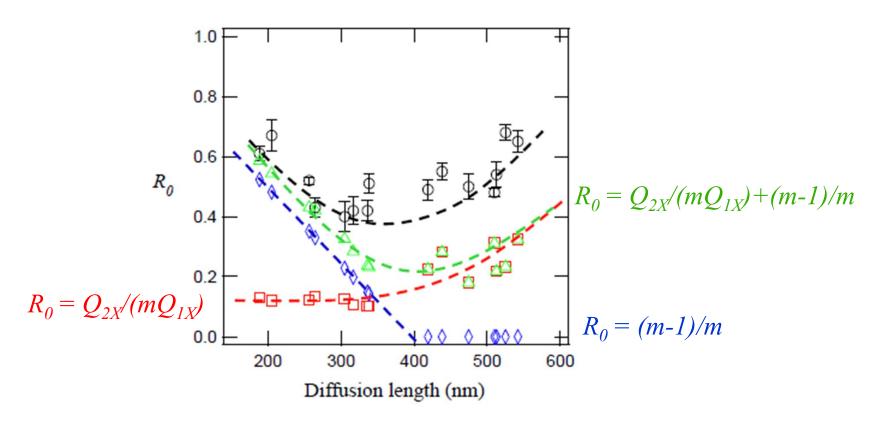
a QD.

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$R_{\theta} \Leftrightarrow ext{Diffusion Length } L_{D}$

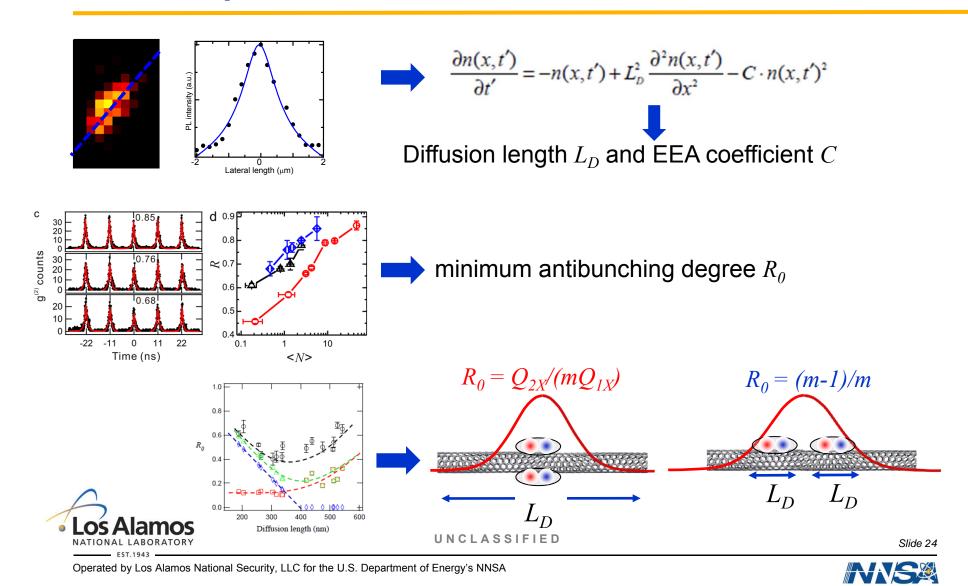




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Summary





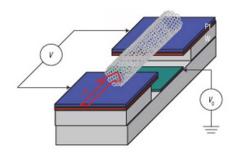
Oxygen Doping Modifies Excitonic Fine Structures of Carbon Nanotube



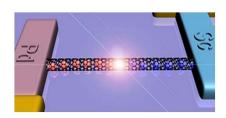
Ma et al. ACS Nano in submission



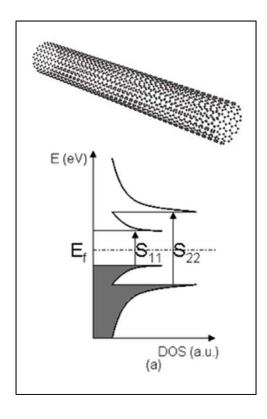
Photoluminescence of SWCNTs: Application

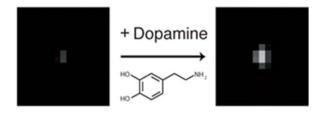


field-effect transistor
Cao et al. Nature Mater. 2005, 4

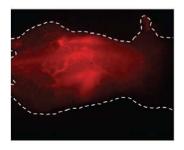


light emitting diode
Wang et al. Nano Lett. 2011, 11





Sensor Kruss et al. J. Am. Chem. Soc. 2014



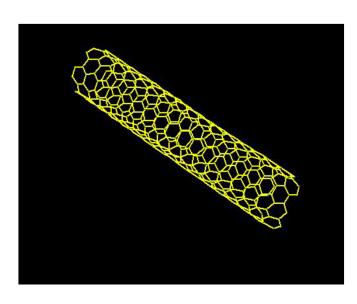
biological imaging
Welsher et al. Nature Nanotech. 2009



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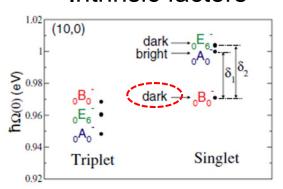
Photoluminescence of SWCNTs: Limitation

External factors

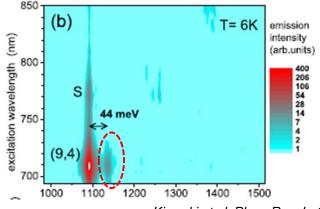


- single layer structure
- structural defects
- adsorbate molecules/atoms

Intrinsic factors



Spataru et al. Phys. Rev. Lett. 2005



Kiowski et al. Phys. Rev. Lett. 2007

dark states below the first bright state



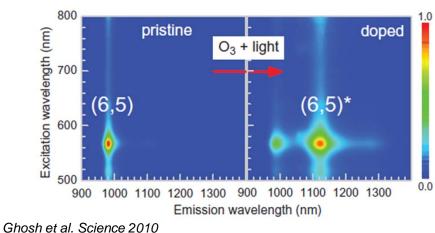
(S. Maruyama)

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Controlled Doping of SWCNTs

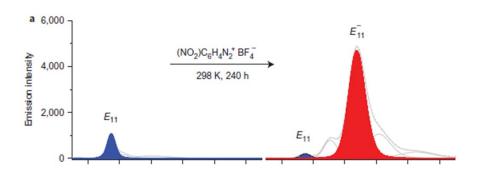
oxygen doping



additional bright, red-shifted peak

diazonium salt doping

Miyauchi et al. Nature Photon. 2013



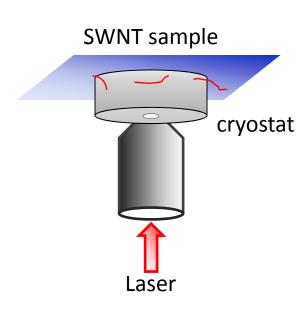
Piao et al. Nature Chem. 2013

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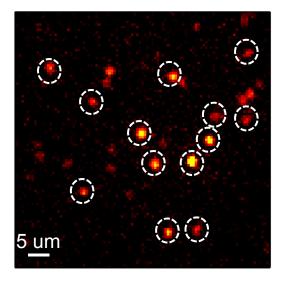




Individual SWCNT Imaging at Low Temperature



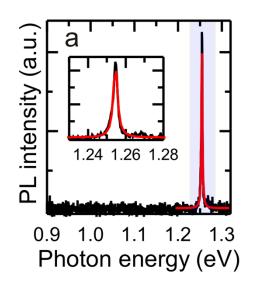
Wide field image

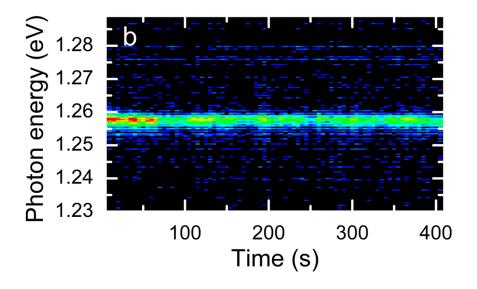


diffraction-limited localized emission



Individual Pristine SWCNT

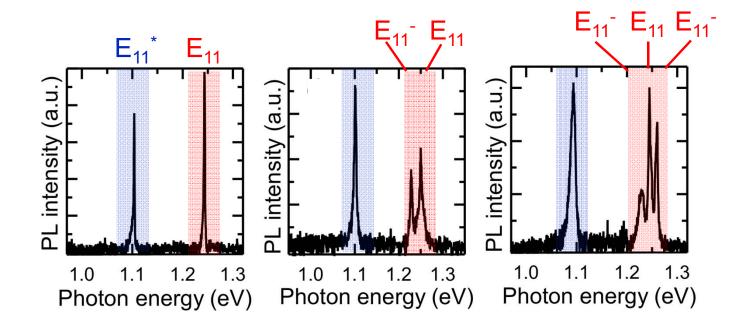




- individual symmetric (88%) peak at ~ 1.25 eV → E₁₁
- linewidth: 4.4 meV



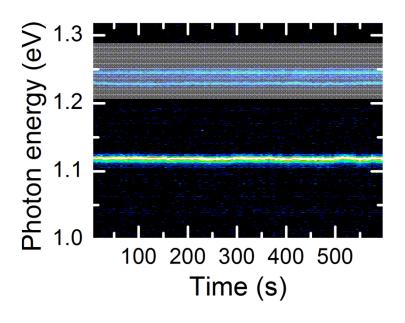
Individual Oxygen-Doped SWCNT

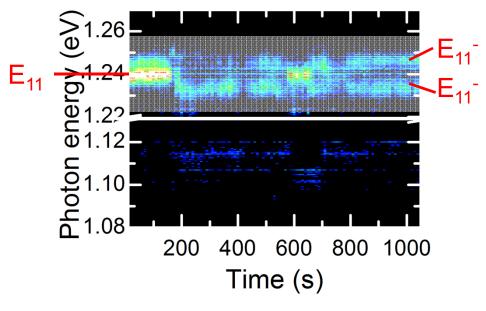


- spectral splitting of E₁₁ peak → 2 to 3 peaks (70%)
- additional red-shifted peak → E₁₁*



Time Evolution of E₁₁ peak





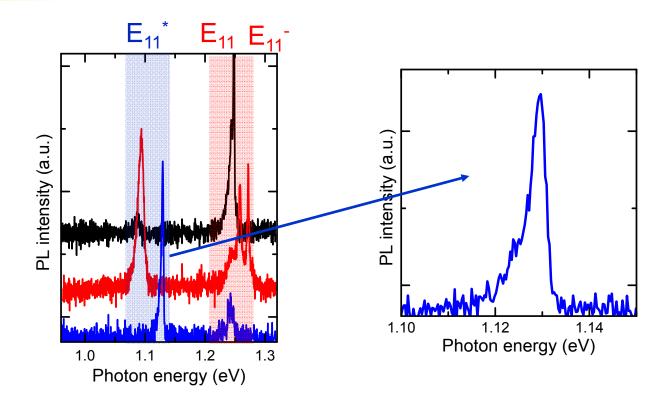
stable peak position and intensity

spectral splitting





Line Shape of E₁₁* Peak

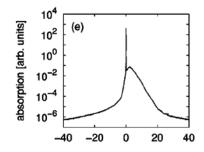


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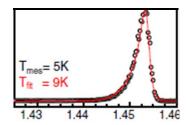
asymmetric peak with long tail at the low energy side (83%)



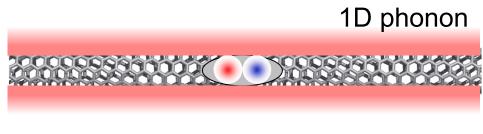
Line Shape of E₁₁* Peak

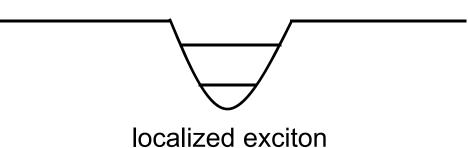


Krummheuer et al. Phys. Rev. B 2002



Galland et al. Phys. Rev. Lett. 2008





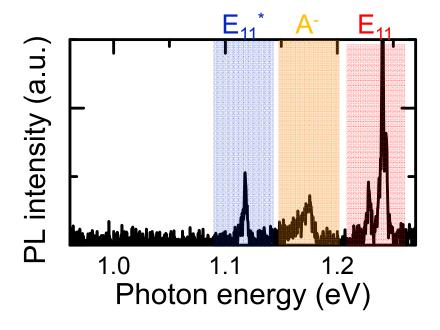
- due to coupling between localized exciton and 1D phonon
- E₁₁* exciton is highly localized.

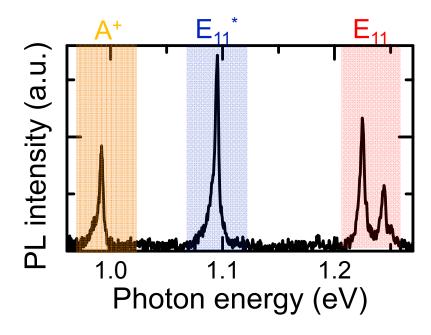


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Additional Peaks

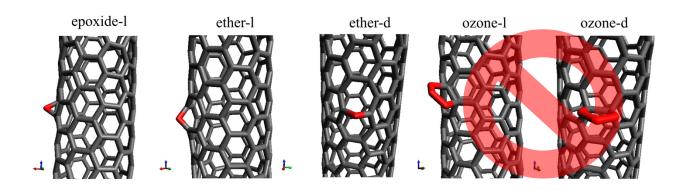




- Both A⁺ and A⁻ are asymmetric with long tail at the low energy side.
- A⁺ and A⁻ excitons are also localized.



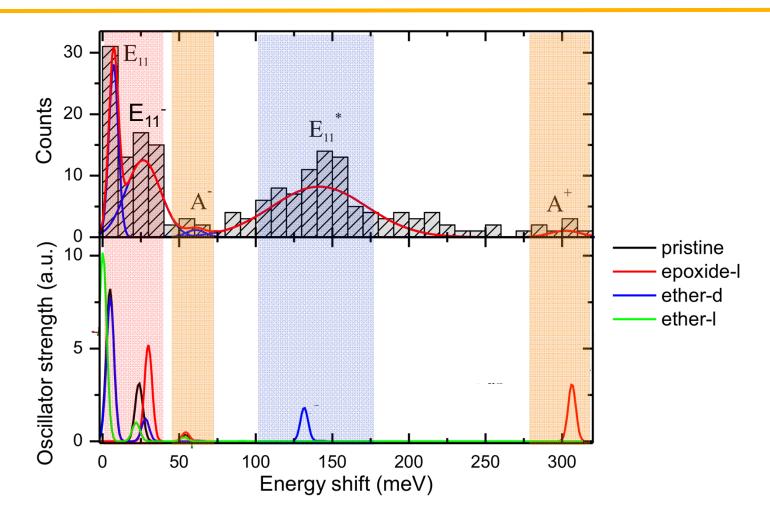
Quantum Chemistry Computations



- Length of tube: 8 nm
- I: along tube axis; d: perpendicular to tube axis
- Geometry optimization: semi-empirical method (AM1 Hamiltonian, MOPAC program)
- Oscillator strengths and transition densities: Time-dependent Density Functional Theory (Gaussian 09, STO-3G basis set, B3LYP functional).



Comparison between Experiment and Simulation

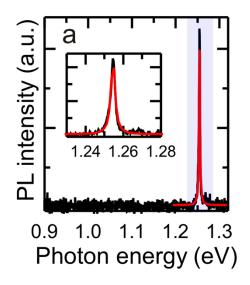


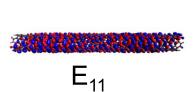


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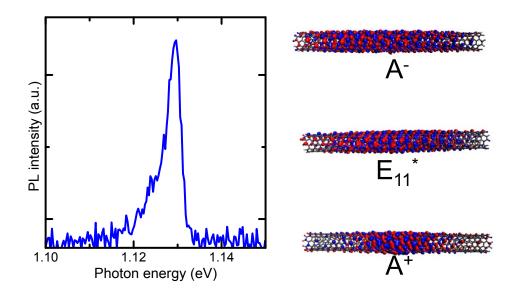


Exciton Transition Density





symmetric lineshape
 → delocalized exciton



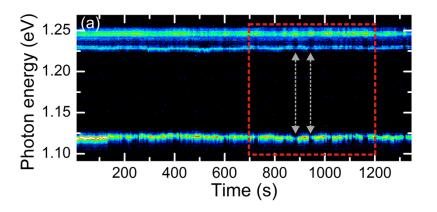
asymmetric lineshape
 → localized exciton

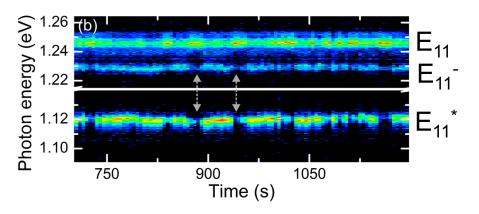


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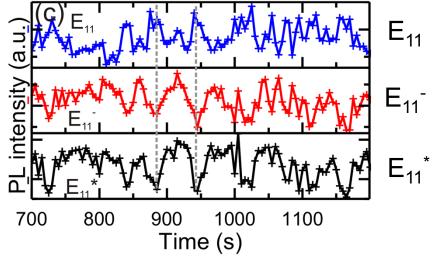


Intensity Correlation between E₁₁ and E₁₁* peaks





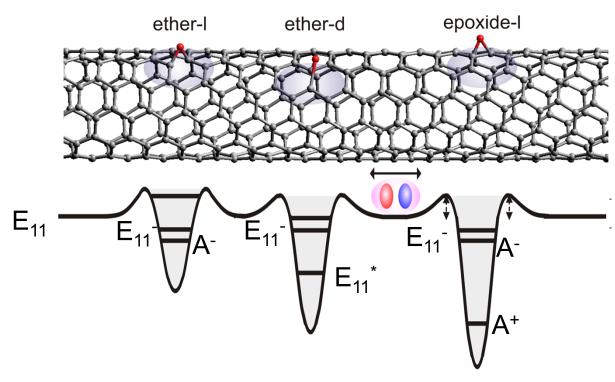
 PL intensities of E₁₁ and E₁₁ peaks are negatively correlated.





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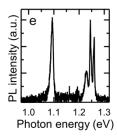
Excitonic Fine Structures



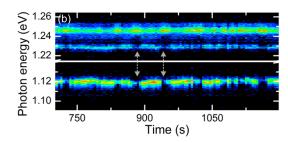
- Oxygen doping introduces potential wells localizing excitons.
- Potential fluctuation leads to negative intensity correlation between E_{11} and E_{11}^* peaks.



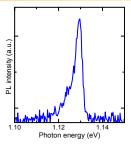
Summary



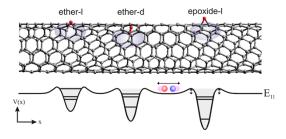
spectral splitting of E₁₁ peak



 negative intensity correlation between E₁₁ and E₁₁* peaks



• asymmetric E₁₁*, A+, and A- peaks



 excitonic fine structures of oxygen-doped SWCNTs



Influence of Exciton Dimensionality on Spectral Diffusion of Single-Walled Carbon Nanotubes



Ma et al. ACS Nano in submission

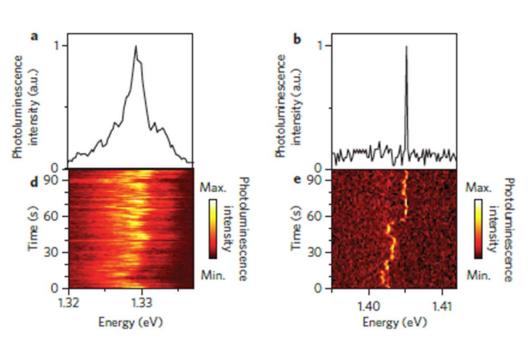
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Spectral Diffusion of SWCNTs

Ideal spectra

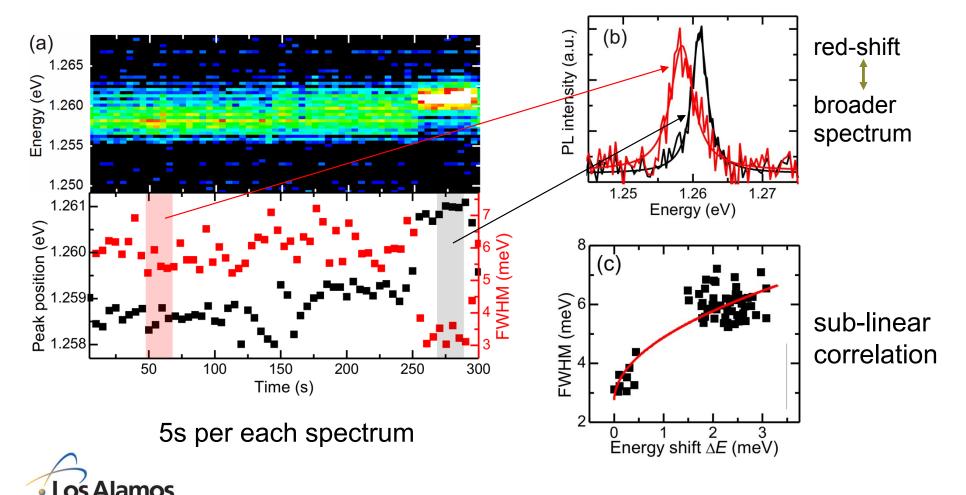
Photoluminescence Time (s) Itime (s) Min. Photoluminescence intensity Min. 1.395 1.405 Energy (eV)

Real spectra



 Quantum information processing requires spectral-diffusion free emission → physical mechanism behind spectral diffusion

PL spectra of Individual SWCNTs at 4K



Quantum-confined Stark effect (QCSE)

CdSe QDs (Bawendi, Science 97) (Bawendi, Science 97) (Bawendi, Science 97) (CdSe QDs (Bawendi,

CdSe/CdS nanorods (Weller, PRL 2004)



2.00 2.02 2.04 2.06 2.08

CdSe/ZnS nanorods (Banin, Nano Lett. 2005)

2.07



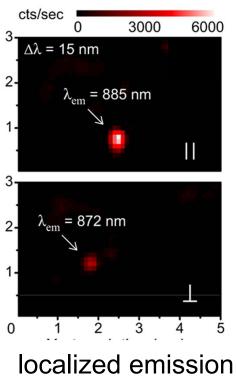


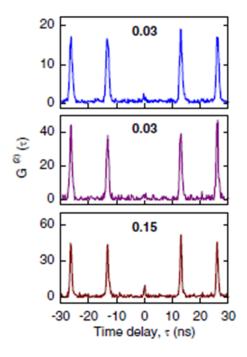
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SWCNTs at Cryogenic Temperature

4 K





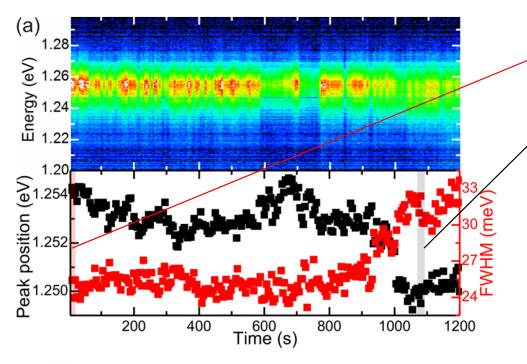
alized emission complete antibunching

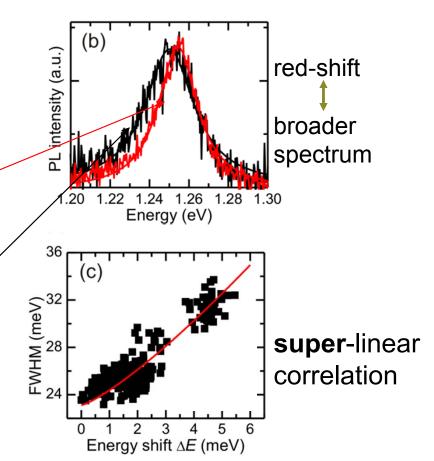
Hoegele et al. Phys. Rev. Lett. 2008

Excitons in SWCNTs at 4K is localized to QD-like excitonic states.
 Sublinear correlation between linewidth and peak shift.

PL spectra of Individual SWCNTs at RT

Exciton in SWCNT is **diffusive** at RT. Can we still observe QCSE?

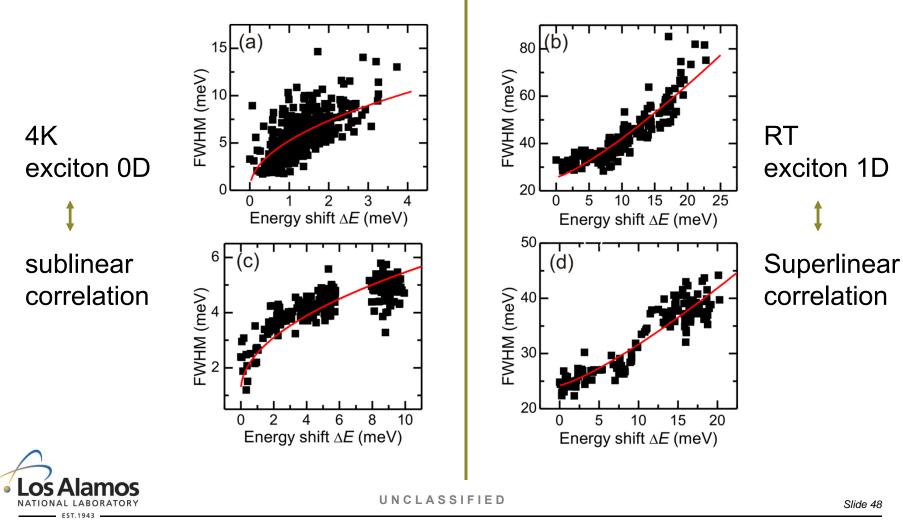




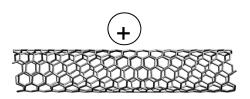


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Exciton Dimensionality ←→ Correlation Coefficient



Single Point Charge Model



For a SWCNT at low electric field, an energy shift of the binding energy (ΔE) caused by QCSE can be described by: $e: electron charge \quad d: tube diameter$

$$\Delta E = k_{\rm b} \frac{\left(edF\right)^2}{E}$$

E: tube binding energy at zero field

 \overline{F} : electric field induced by surface charges

 $k_{\rm b}$: constant

If we use a Gaussian stochastic model to describe spectral diffusion behavior:

$$\langle \delta(\Delta E)(t)\delta(\Delta E)(0)\rangle = \sigma^2 e^{-t/t_c}$$

 σ : spectral broadening in meV

At time t=0 we obtain

$$\langle \delta^2(\Delta E)(0) \rangle = \sigma^2$$

In a simple assumption, we assume that there is an external point charge located on the surface of a SWCNT that creates a local electric field:

$$F = \frac{q}{4\pi\varepsilon r^2}$$

 ε : permittivity of the surrounding environment

r: exciton-charge distance



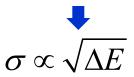
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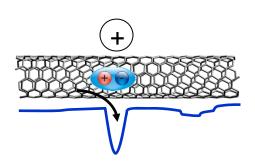
Single Point Charge Model

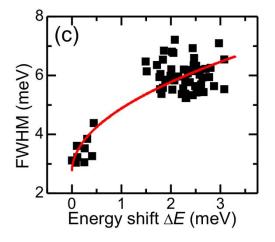
At cryogenic temperature:

- Excitons are localized to quantumdot-like states.
- Exciton-charge distance r and local environment permittivity ε: constants.
- Exciton binding energy fluctuation is mainly induced by fluctuations in surface charges:

$$\delta F = \frac{1}{4\pi\varepsilon r^2} \delta q$$







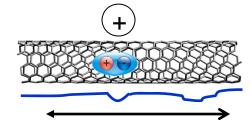


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Single Point Charge Model

At room temperature:

- Excitons are diffusive (~100 nm).
- Large fluctuation in exciton-charge distance r.

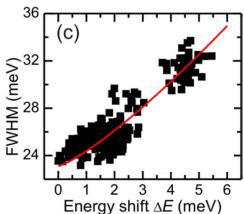


 Contributions of surface charge/dielectric environment fluctuation become negligible.

$$\delta F = -\frac{2q}{4\pi\varepsilon r^3} \delta r$$



$$\sigma \propto \Delta E^{5/4}$$

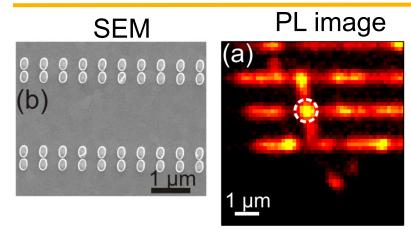




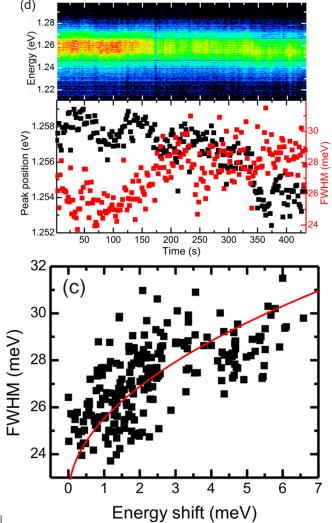
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SWCNTs Coupled to Surface Plasmons



 $\sigma \propto \sqrt{\Delta E}$ \blacksquare exciton "localization" at room temperature

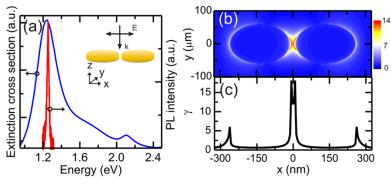




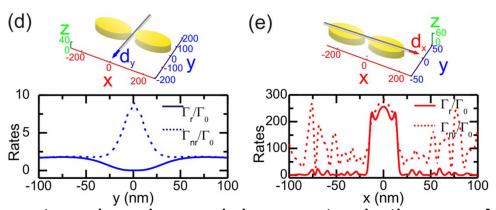
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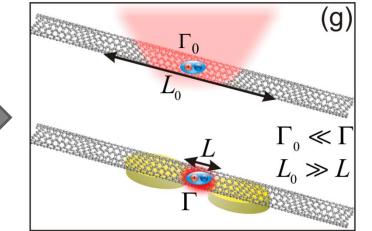
Numerical Simulations



highly confined excitation regine (<30nm)



strongly enhanced decay rates in the gap -> chance of exciton diffusion greatly decreased.

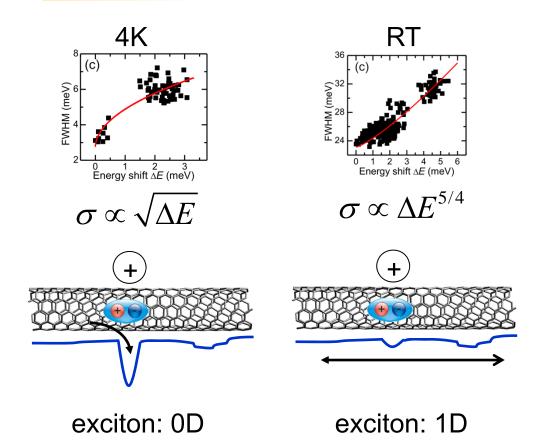


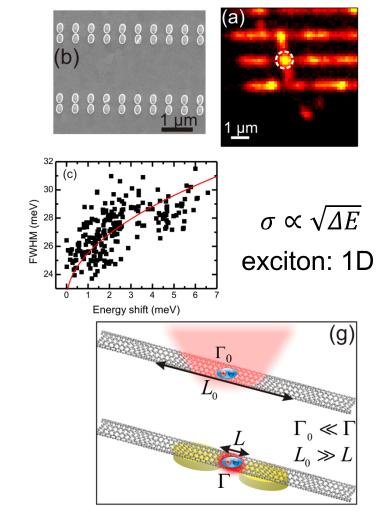






Summary





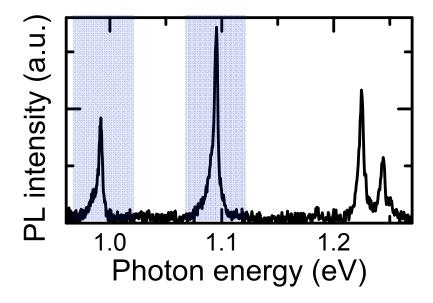


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Current Project

Realization of Telecom Range (1300 - 1550 nm) Single Photon Source from SWCNTs at Room Temperature

- Air stable oxygen-doped SWCNTs
- Photon statistics of 1100 nm and 1300 nm peaks





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Thanks for your attention!

